



Departures from gradient wind balance during secondary eyewall formation

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Background

It has been proposed that the **essential dynamics of secondary eyewall formation (SEF)** are related to departures of (gradient wind) balance in the boundary layer (Huang et al. 2012; Abarca and Montgomery 2013). In this view there is a progressive boundary layer control that includes the development of supergradient flow. The latter is crucial for SEF in:

- The establishment of the sustained deep convection
Result of the strong horizontal convergence (a gradient force pointing outward in a region of inflow)
- The emergence of the wind maximum
Result from mean influx of absolute vorticity surpassing the depletion of absolute angular momentum

To what extent do full-physics numerical simulations undergoing SEF depart from axisymmetric balance evolution as defined by Shapiro and Willoughby (1982)?

Methodology

Compare:

- Radial velocity
- Vertical velocity
- Tangential velocity tendency

RAMS Vs Sawyer-Eliassen

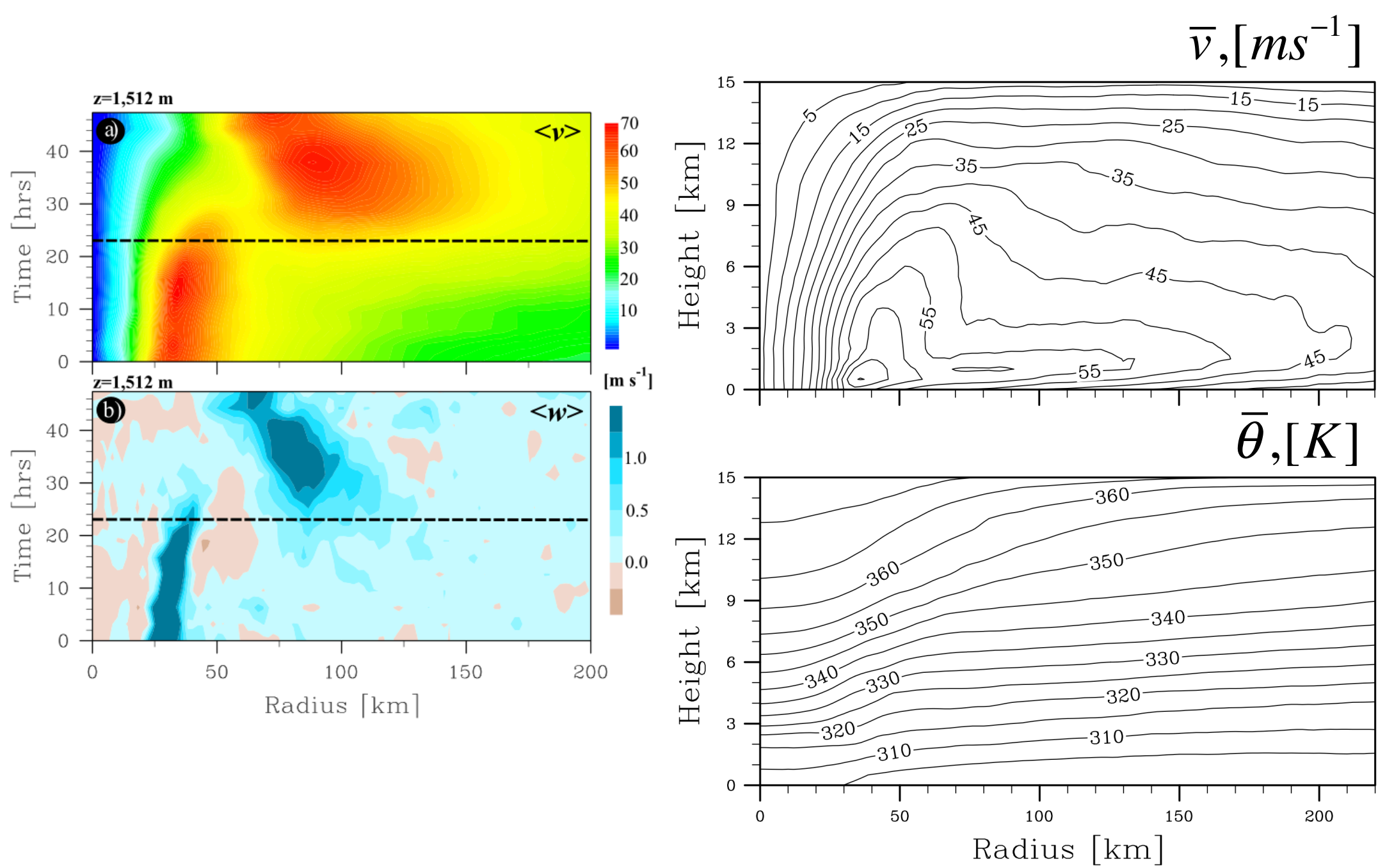
Sawyer-Eliassen

$$\frac{\partial}{\partial r} \left[-g \frac{\partial \chi}{\partial z} \frac{1}{pr} \frac{\partial \psi}{\partial r} - \frac{\partial(\chi C)}{\partial z} \frac{1}{pr} \frac{\partial \psi}{\partial z} \right] + \frac{\partial}{\partial z} \left[\left(\xi \chi \zeta_a + C \frac{\partial \chi}{\partial r} \right) \frac{1}{pr} \frac{\partial \psi}{\partial z} - \frac{\partial(\chi C)}{\partial z} \frac{1}{pr} \frac{\partial \psi}{\partial r} \right] = g \frac{\partial(\chi^2 Q)}{\partial r} + \frac{\partial(C \chi^2 Q)}{\partial z} - \frac{\partial(\chi^2 \xi F_\lambda)}{\partial z}$$

- Integrated using the successive overrelaxation scheme (Press *et al.* 1992)

RAMS

- Full-physics, mesoscale integration
- Exhibited a canonical SE
- Described in Terwey and Montgomery (2008) and Abarca and Montgomery (2013)



Ellipticity condition

Met when $D < 0$, where:

$$D = B^2 - 4AC \quad \chi = \frac{1}{\theta}$$

$$A = -g \frac{\partial \chi}{\partial z}$$

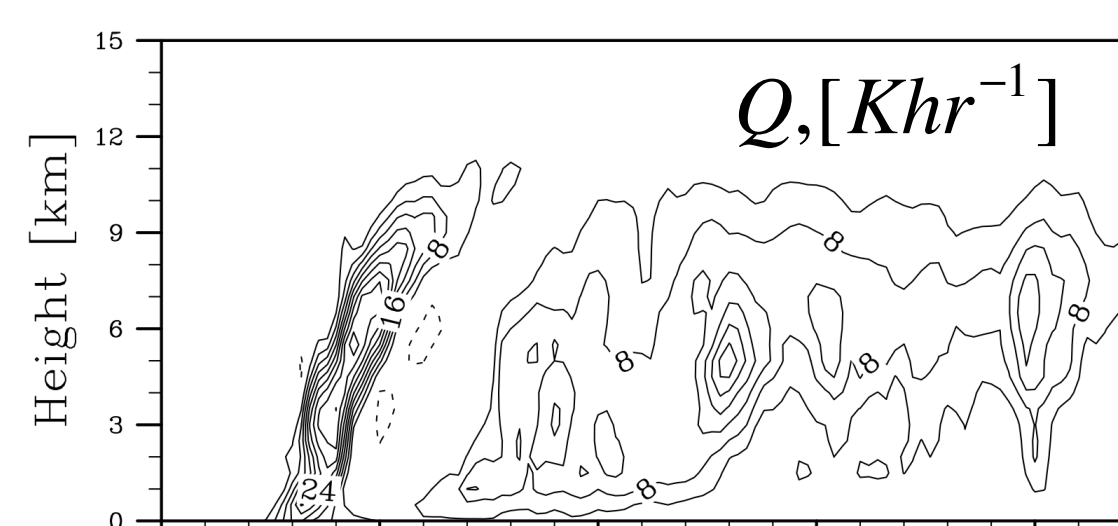
$$B = \frac{\partial \left[\chi \left(\frac{v^2}{r} + f_c v \right) \right]}{\partial z}$$

$$C = \left(\frac{2v}{r} + f_c \right) \chi \zeta_a + \left(\frac{v^2}{r} + f_c v \right) \frac{\partial \chi}{\partial r}$$

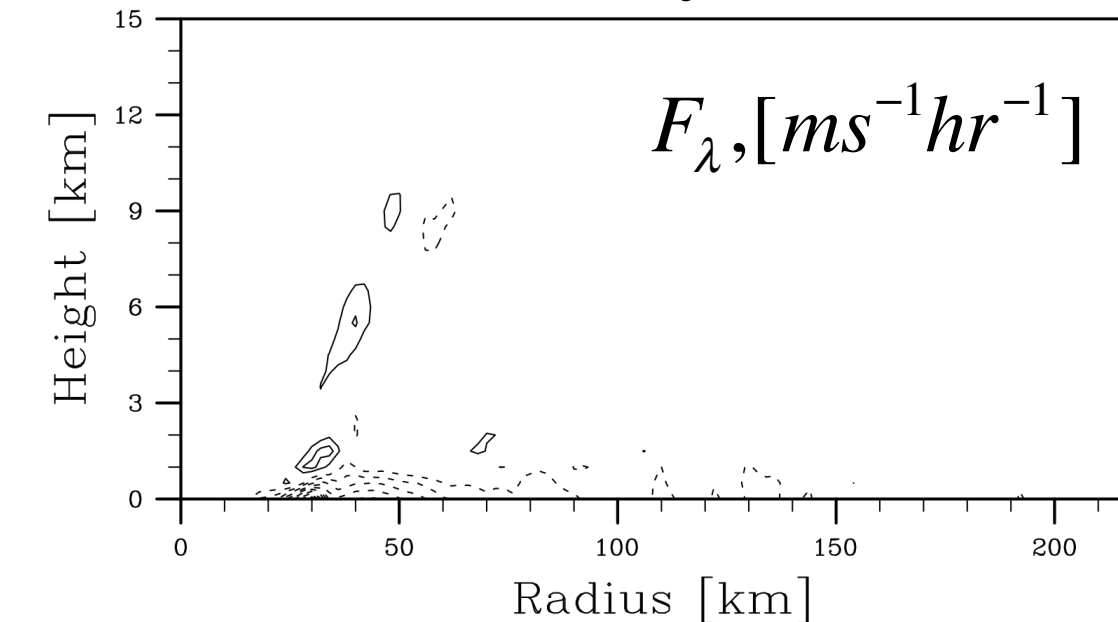
- Regularization procedure

Forcing Functions

$$Q = \bar{\theta} - u' \frac{\partial \theta'}{\partial r} - \frac{v'}{r} \frac{\partial \theta'}{\partial \lambda} - w' \frac{\partial \theta'}{\partial z}$$



$$F_\lambda = \frac{\partial \bar{v}}{\partial t} + \bar{u} \zeta_a + \bar{w} \frac{\partial \bar{v}}{\partial z} \quad (\text{by residual})$$



Definition of balance

Three options:

a) Azimuthal averages

As performed before (e.g. Persing and Montgomery 2003), but not in balance

a) Smith et al. (2006)

Use azimuthal average tangential wind field and derive the temperature field in balance with it

a) Holton (2004)

Use azimuthal average temperature field and derive the tangential wind field in balance with it

Take home message:

Axisymmetric balanced dynamics of Shapiro and Willoughby (1982) are unable to capture SEF spin-up

RAMS

Balanced fields

D

u

w

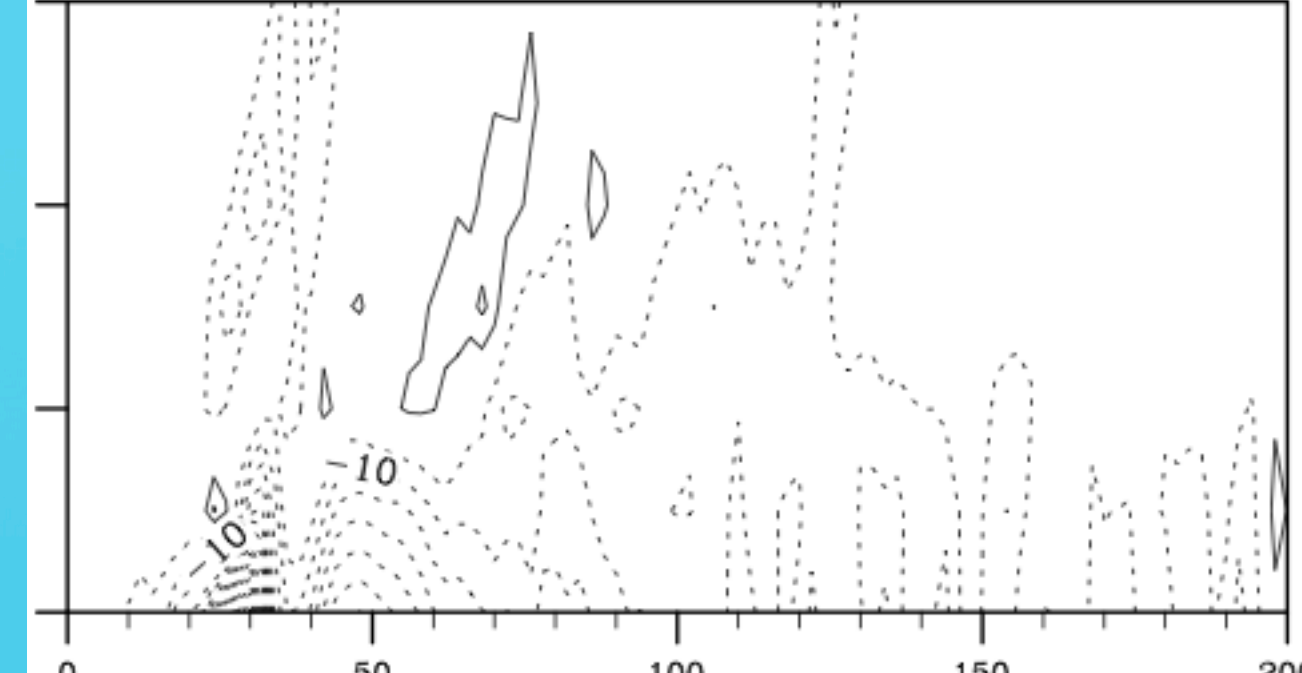
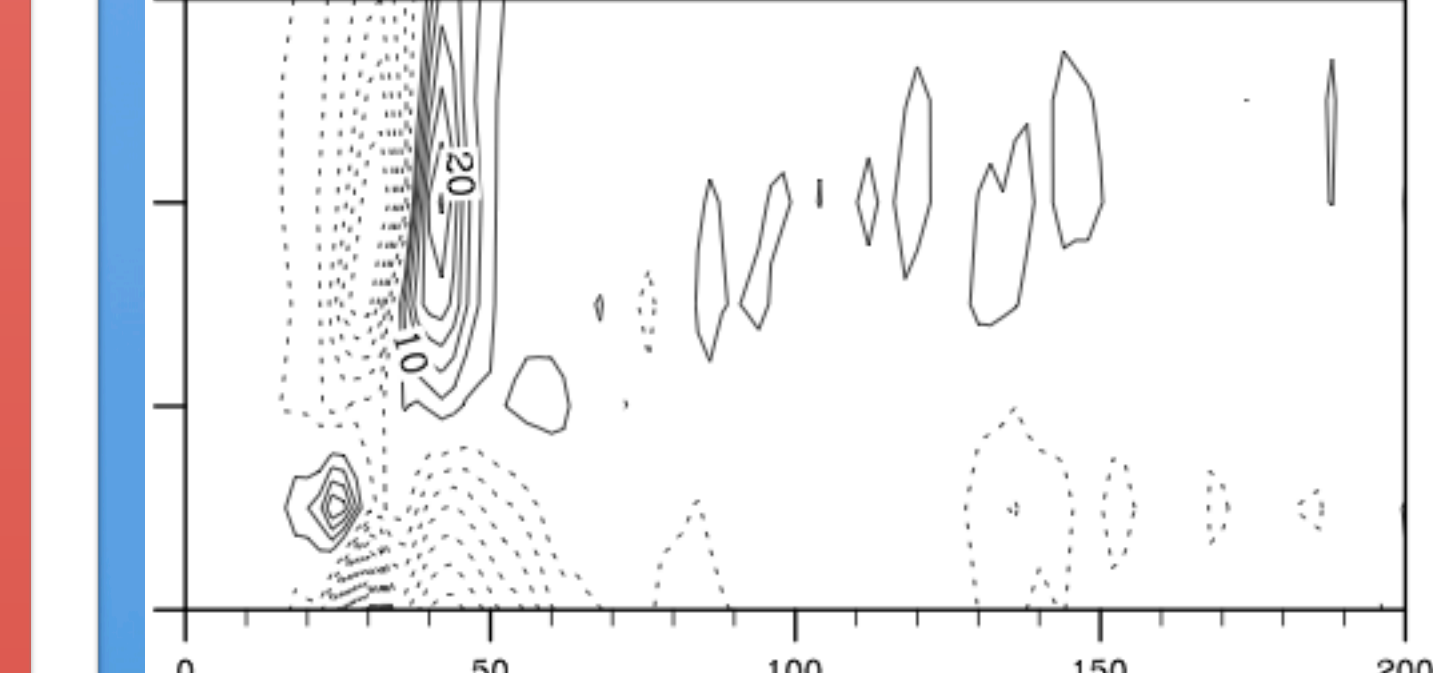
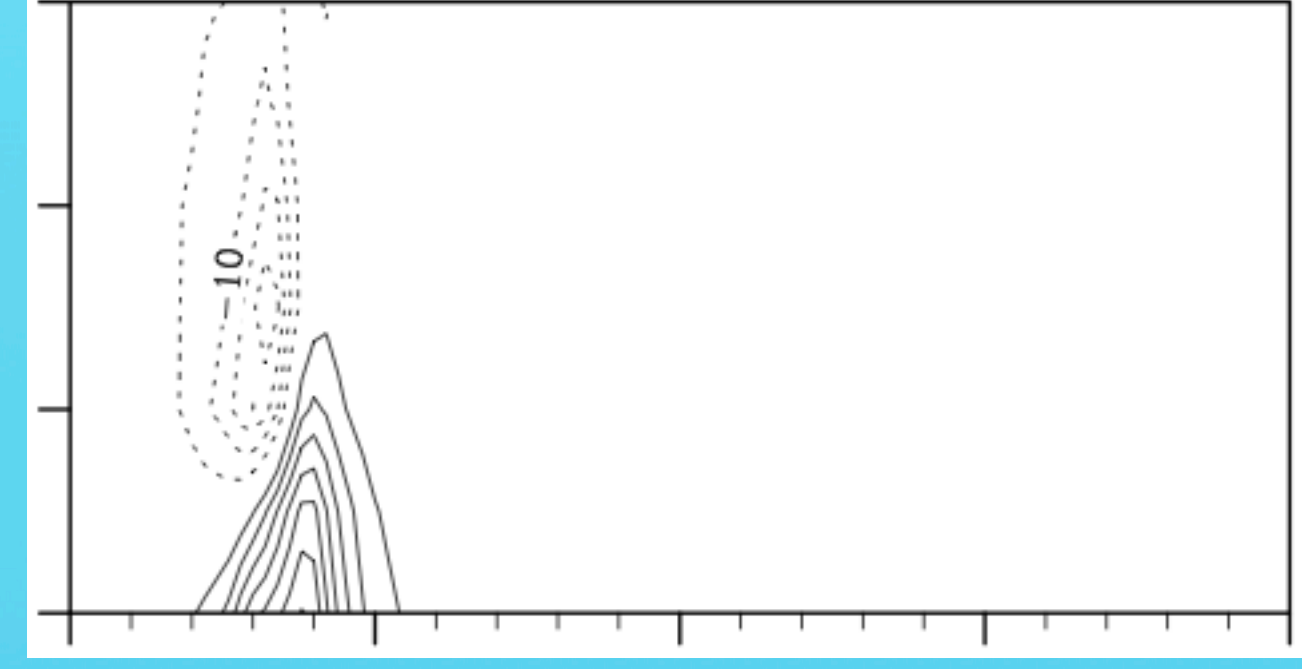
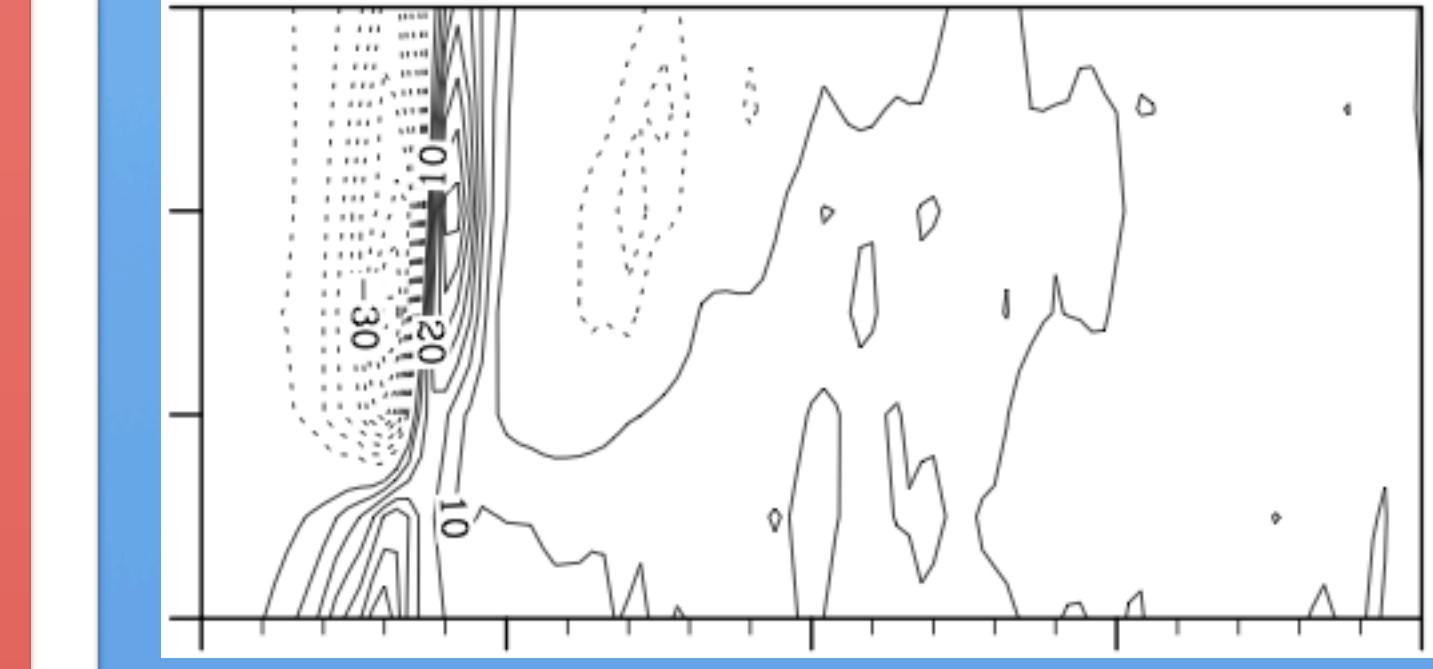
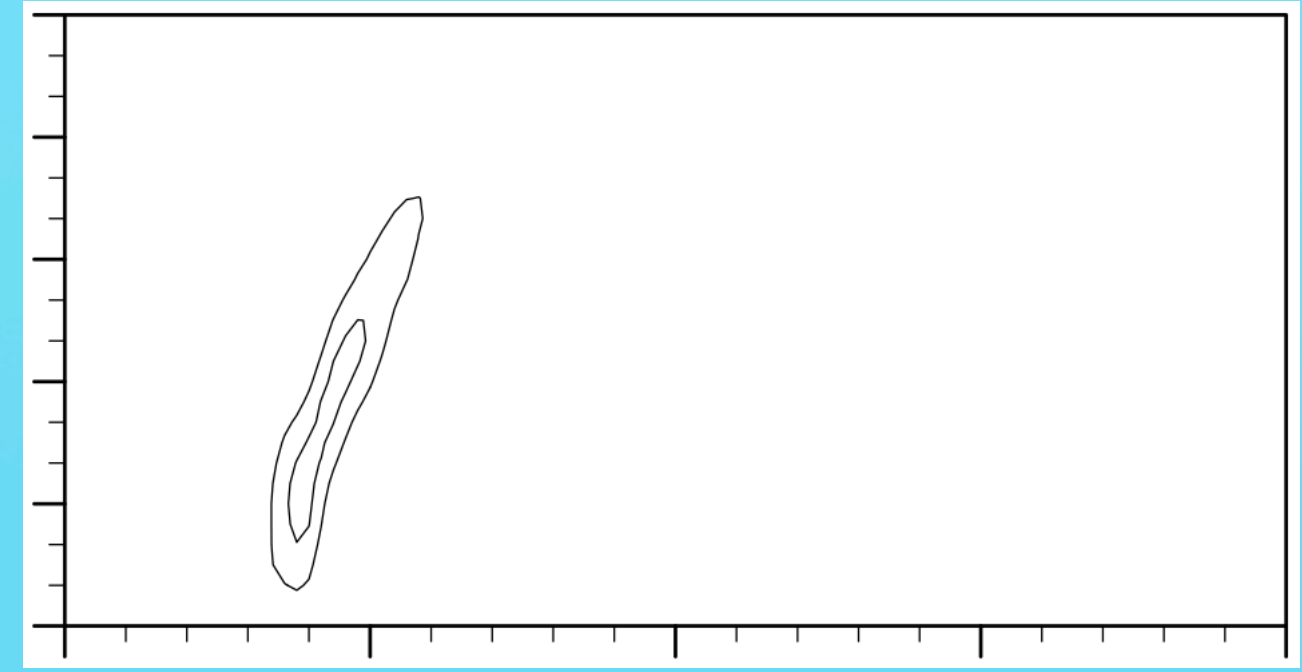
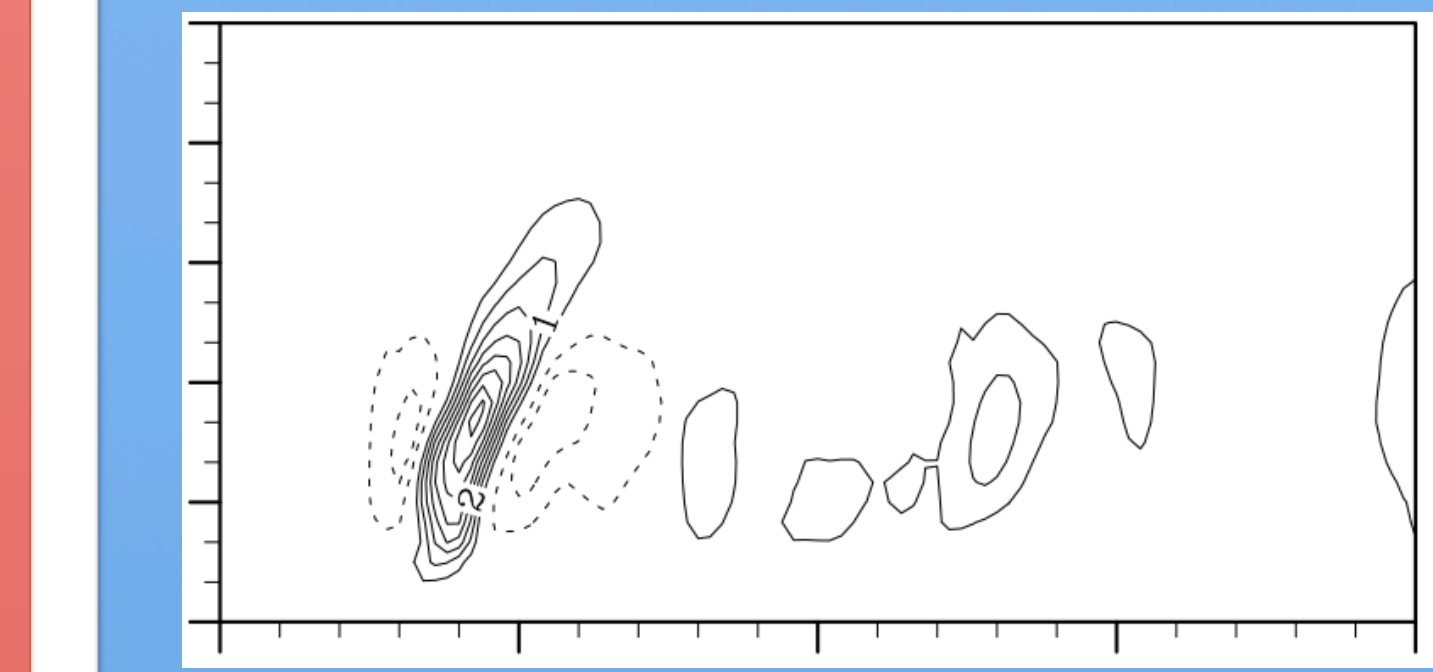
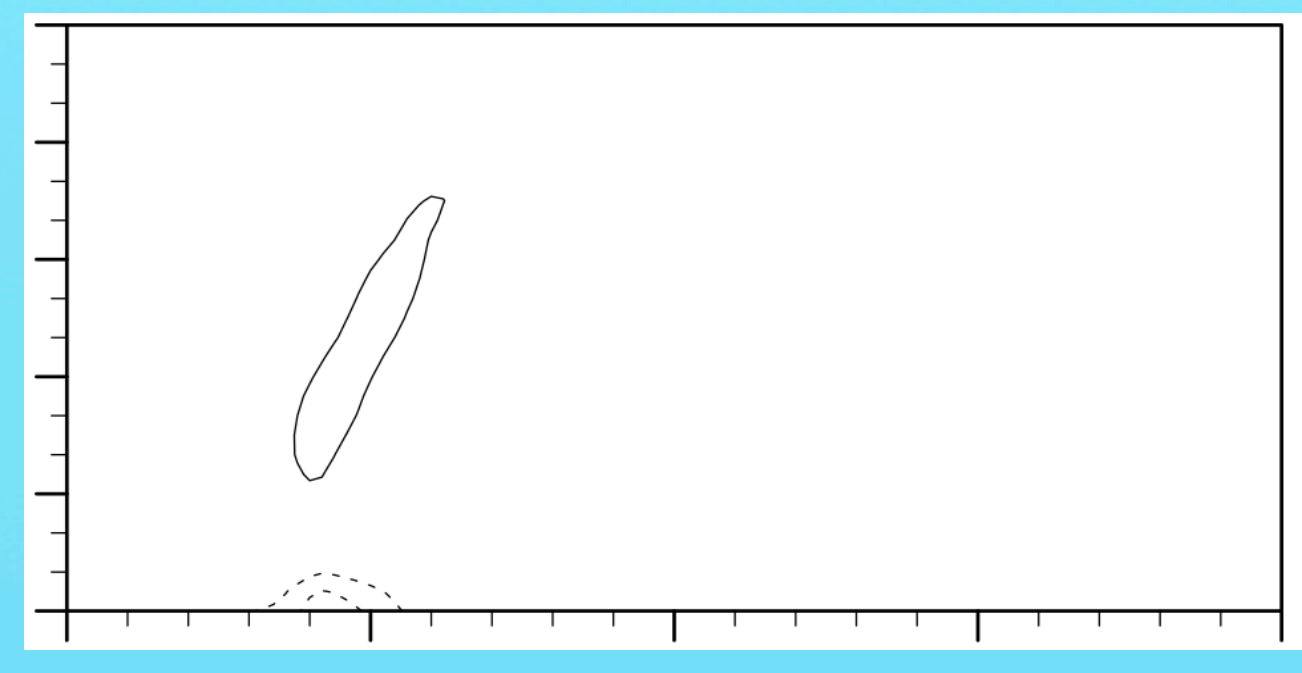
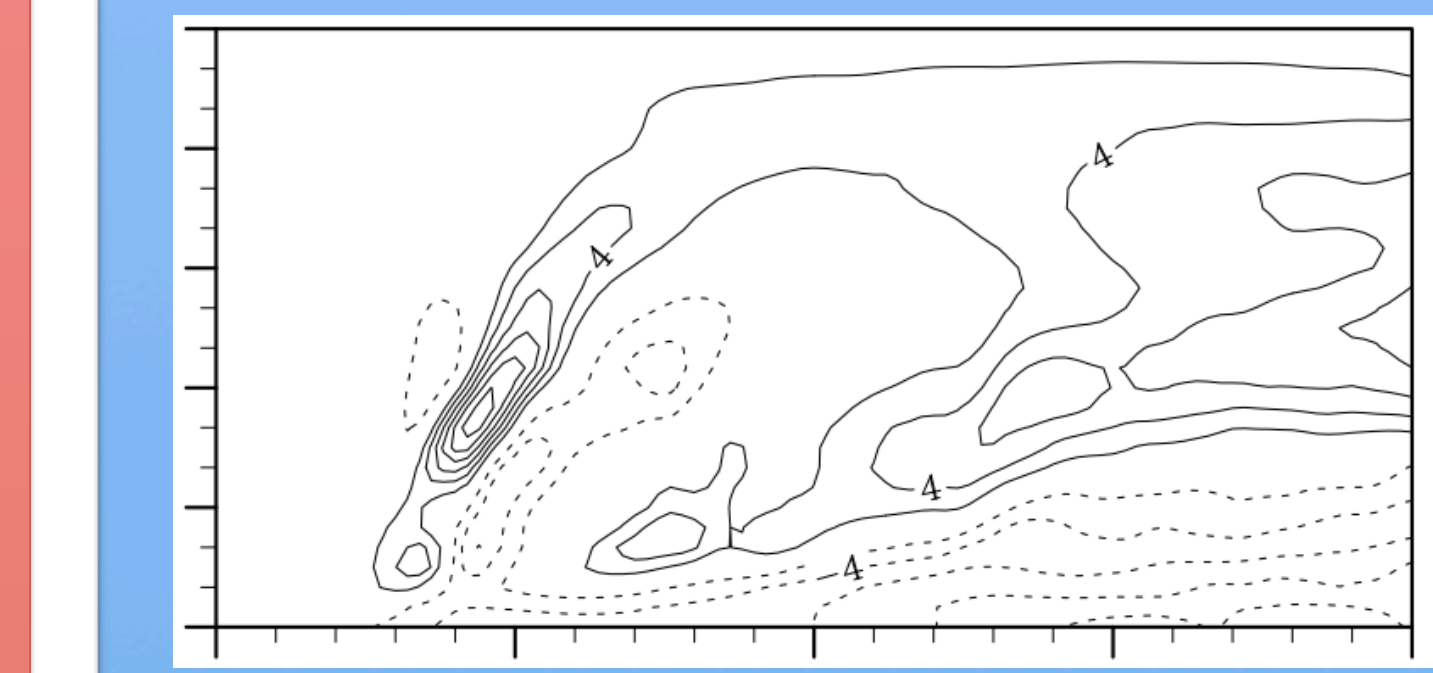
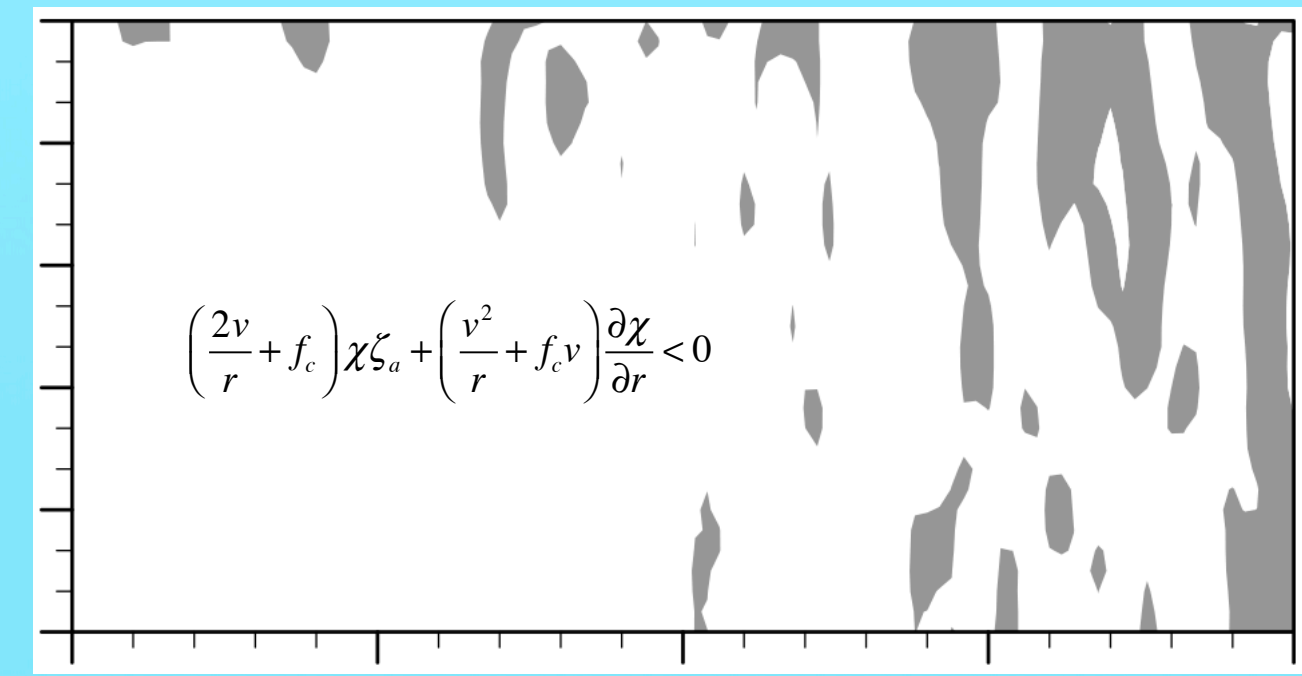
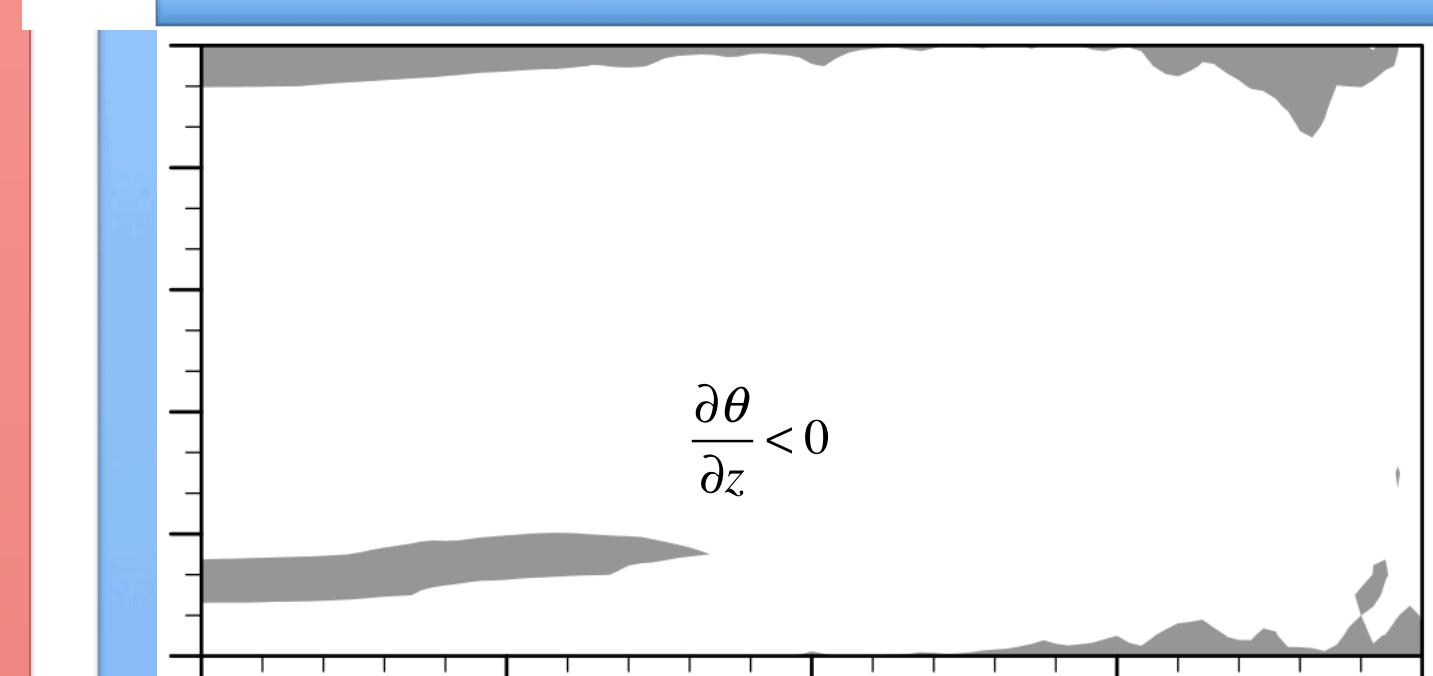
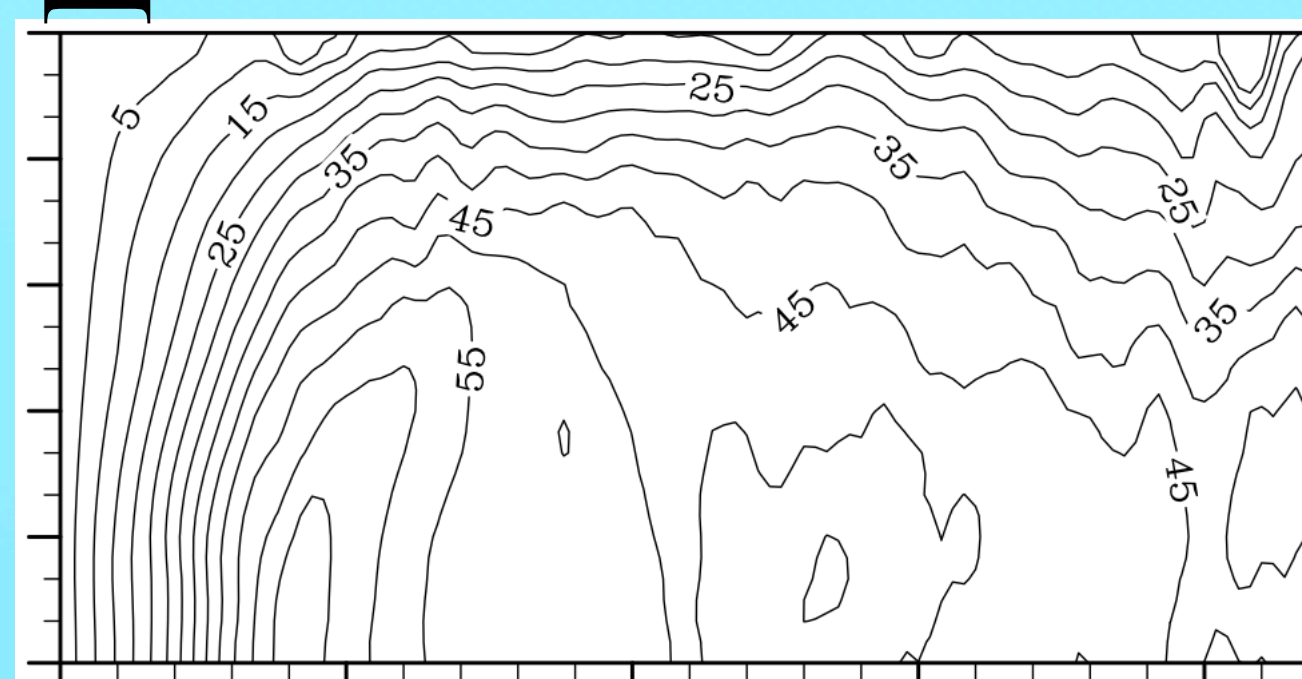
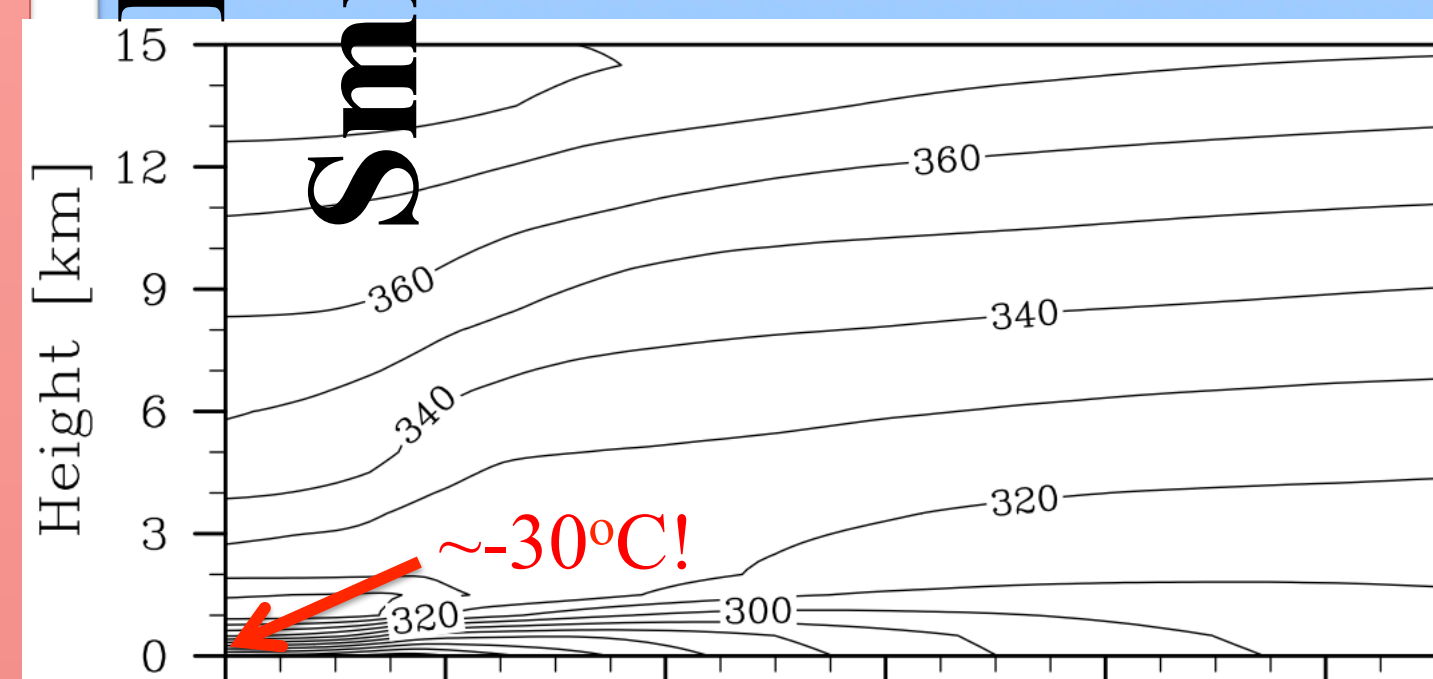
$$\frac{\partial v}{\partial t} = -u \zeta_a - w \frac{\partial v}{\partial z}$$

$$\frac{\partial v}{\partial t} = -u \zeta_a - w \frac{\partial v}{\partial z} + F_\lambda$$

Azimuthal averages

Balanced fields
Smith et al. (2006)

Balanced fields
Holton (2004)



References

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